Room-temperature operation of a p-type molecular spin photovoltaic device on a

transparent substrate

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Abstract:

The coupling of diverse degrees of freedom opens the door to physical effects that go beyond

each of them individually, making multifunctionality a much sought after attribute for high

performance devices. Here, we show the multifunctional operation of a single layer p-type

organic device, displaying both spin transport and photovoltaic effect at the room temperature

on a transparent substrate. The generated photovoltage is almost three times larger than the

applied bias to the device which facilitates the modulation of the magnetic response of the

device with both bias and light. The device shows an increase in power conversion efficiency

under magnetic field, an ability to invert the current with magnetic field and under certain

conditions it can act as a spin-photodetector with zero power consumption in the standby mode.

The room temperature exploitation of the interplay among light, bias and magnetic field in the

single device with a p-type molecule opens a way towards the development of efficient high-

performance spin photovoltaic cells.

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Vertical spin valve (SV) structures are widely used to study spin transport through organic semiconductors (OSCs) [1-3]. Unlike inorganic semiconductors, OSCs are composed of light chemical elements resulting in a weak spin-orbit coupling and hyperfine interaction, which leads to long spin diffusion times^[4-8]. Recently, and in spite of the relatively low carrier mobility of the materials studied, long distance spin transport has been achieved in several organic semiconductor based spin valves^[3,9]. OSCs also offer a large degree of tunability in their mechanical, electrical and optical properties over their inorganic counterparts, which can be used to add multifunctionality in a single device^[10-17]. One such example is the observation of a photovoltaic effect in an electron transporting (n-type) C₆₀ based spin valve, in which the magnetic and optical response are coupled and this leads to several novel functionalities in a single device at low temperature^[16]. In spite of this advancement, a basic standard small molecule organic photovoltaic cell is composed of both an electron (n-type) and a hole transporting (p-type) layer^[18–20]. In order to achieve a high photovoltaic response of the SV towards its room temperature multifunctional operation, devices with a hole transporting material have to be realized as a necessary step.

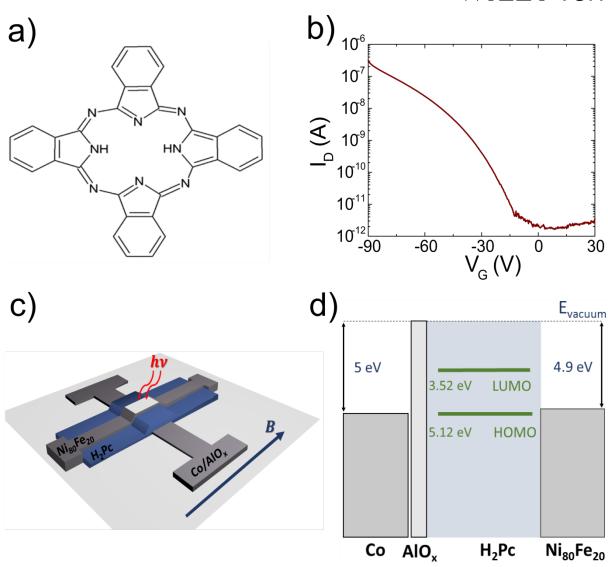


Figure 1. a) Schematic of a hydrogen phthalocyanine (H_2Pc) molecule. b) Transfer characteristic of a 30nm thick H_2Pc based lateral organic field effect transistor ($L = 10 \mu m$; W = 1 mm) showing p-type operation with a source-drain bias $V_D = 40V$. c) Schematic presentation of the H_2Pc based spin-valve operating as a spin-photovoltaic device on a transparent glass substrate. d) Rigid band energy diagram of the device. The Fermi level (E_F) of both Co and $Ni_{80}Fe_{20}$ electrodes matches very well to the HOMO level of the H_2Pc molecule, favoring a hole dominated transport through the device.

Here in this manuscript we present the room temperature operation of a p-type molecular spin photovoltaic (MSP) device on a transparent glass substrate. Our device has a simple vertical SV geometry composed of a bottom ferromagnetic (FM) electrode cobalt (Co) capped by a leaky AlO_x barrier, a p-type organic molecule (hydrogen phthalocyanine; H₂Pc) acting as the

spin transporting channel and a top FM electrode ($Ni_{80}Fe_{20}$) ^[2,3]. H₂Pc is a commonly used p-type material for small molecule organic photovoltaic cells and is also known for its stability in ambient conditions^[20–22]. **Figure 1a** presents the schematic diagram of an H₂Pc molecule, while Figure 1b shows the transfer characteristic of a 30 nm thick H₂Pc based lateral organic field effect transistor (LOFET) respectively, which indicates the p-type operation with a lateral hole mobility of 3.3×10^{-5} cm²V⁻¹s⁻¹ ^[23]. The growth of the H₂Pc molecule on the glass substrate is highly amorphous thus it exhibits isotropic charge transport and we do not expect large changes in the charge carrier mobility in a vertical geometry (See Figure S1). Figure 1c presents the schematic of the complete vertical SV device and Figure 1d represents the rigid band energy diagram of the device. The Fermi energy (E_F) of both the FM electrodes is well matched with the highest occupied molecular orbital (HOMO) of H₂Pc with a hole injection barrier of 0.22 eV while the energy barrier for the electron injection into the lowest unoccupied molecular orbital (LUMO) is around 1.5 eV ^[24]. Hence, in our SV device, the charge carrier transport is most likely to be dominated by holes.

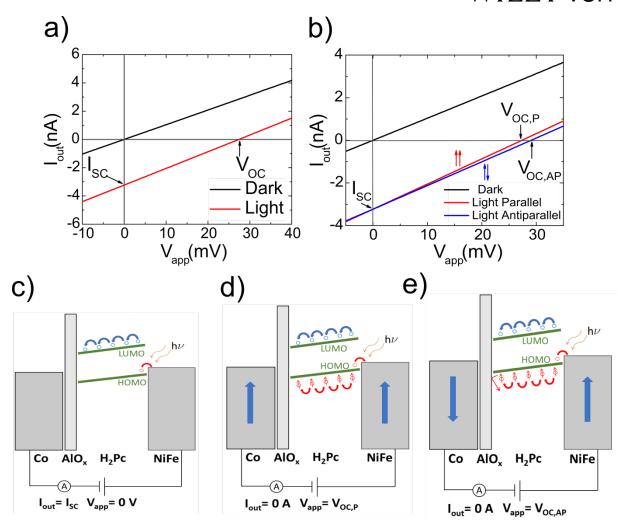


Figure 2. a) Current-voltage (I-V) response of the MSP device in dark and light (7.5 mW.cm²) irradiation conditions without the application of any magnetic field. b) I-V response of the device in dark and during light irradiation when the electrodes are in parallel or anti-parallel configuration. c) Mechanism of electron-hole transport upon light irradiation under zero bias conditions. d)-e) Spin dependent charge carrier transport through the molecular layer under light irradiation in both parallel and anti-parallel configurations of the electrodes in the open circuit mode.

Figure 2a represents the current-voltage (I-V) response of the device, both in dark and light irradiation conditions, when the voltage bias is applied to the top $Ni_{80}Fe_{20}$ electrode and the bottom Co electrode is grounded. The device shows a photovoltaic response with a short circuit current (I_{SC}) of 3.2 nA and an open circuit voltage (V_{OC}) of close to 27 mV at room temperature

for an incident light irradiation of 7.5 mW.cm⁻² when no external magnetic field is applied to the device. The values of I_{SC} and V_{OC} are relatively small as compared to other conventional photovoltaic cells since in our case the device is composed of a single photoactive layer^{[14,15,25–} ^{28]}. Under an applied external magnetic field, the photovoltaic response of the device increases, yielding a higher V_{OC} value when the orientation of the FM electrodes are antiparallel as shown in Figure 2b. The power conversion efficiency increases \approx 12% when the relative orientation of the FM electrodes changes from parallel to antiparallel state [16]. The mechanism of the photovoltaic effect under the absence of applied magnetic field and for parallel and anti-parallel orientation of the electrodes under magnetic field is schematized in Figure 2c-2d in the rigid band approximation at equilibrium. We infer from the I-V response of our device that, at equilibrium, there is an interface dipole formed at the H₂Pc/ Ni₈₀Fe₂₀ interface despite of the apparently good energy level alignment. On the contrary, the bottom interface is less affected because of the presence of the AlO_x layer which isolates the molecule from highly reactive Co electrode underneath. In our model (Figure 2c), the photo-generated holes at the molecule-Ni₈₀Fe₂₀ interface are collected by the Ni₈₀Fe₂₀ electrode due to the in-built potential, before recombining with the electrons. The electrons then can either be transported to the bottom Co electrode by the in-built potential or lost due to partial recombination. Since H₂Pc is a p-type organic semiconductor, the generated photocurrent is mainly dominated by the collected holes at the molecule-Ni₈₀Fe₂₀ interface. In the open-circuit mode, the injected spin-polarized carriers must compensate the photogenerated carriers or in other words, the applied bias must be equal and opposite to the photo-generated voltage. In the parallel configuration of the electrodes, the photo-generated voltage (V_{OC, P}) is the same as in the case without applied magnetic field while in the anti-parallel case, the injected spin-polarized holes are reflected by the bottom Co electrodes and can only be compensated by an enhanced photo-voltage Voc, AP (Figure 2d and 2e).

We define this change in open-circuit voltage (ΔV_{OC}) as $V_{OC, AP} - V_{OC, P}$. Figure 3a presents the modulation of ΔV_{OC} with respect to magnetic field and light. At stronger light intensities, the photo-generated carriers can be compensated with the application of a higher device bias in the parallel configuration and hence, even a higher applied bias is needed for the compensation in the antiparallel configuration. We observe that the value of ΔV_{OC} increases linearly with increasing light intensity (Figure 3b). This can be useful for a spin-photodetector with zero power consumption since in the open-circuit mode of operation the current through the device is zero [29–31]. In the short circuit-mode, the photogenerated current remains constant over the range of applied magnetic field simply because the photogenerated carriers are not spin-polarized (Figure 2c).

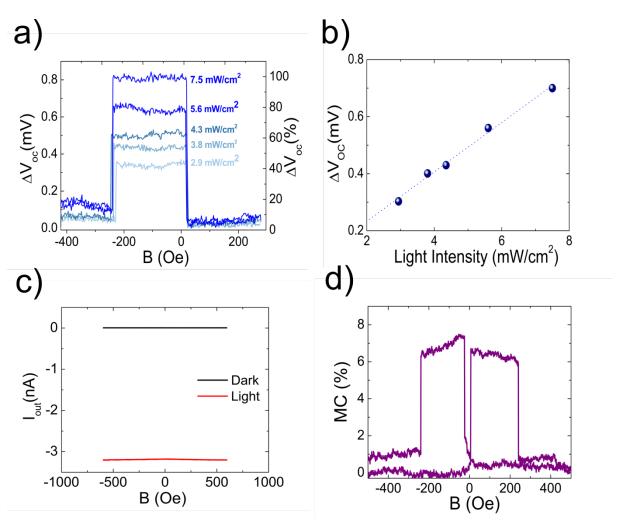


Figure 3. a) Modulation of the open-circuit voltage with magnetic field at various light intensities. The net current through the device is zero. b) The difference in open-circuit voltage $(\Delta V_{OC} = V_{OC, AP} - V_{OC, P})$ as a function of light intensities. ΔV_{OC} is linear with respect to the light intensities. c) I-B response of the device with zero applied bias in dark and light irradiation conditions. The photo-generated current is unaffected on the application of magnetic field as it is not spin polarized. d) Magnetocurrent (MC) response of the device at 300K.

In dark conditions, we define the magnetocurrent (MC) (in percentage) as $(I_P - I_{AP})/I_{AP} \times 100 \%$, where I_P and I_{AP} are the currents in the parallel and anti-parallel configuration of the electrodes respectively. We observe an MC of 7% for an applied bias of 10 mV at room temperature (Figure 2d). Since the photo-generated voltage (of around 27 mV) with a light irradiation of 7.5 mW.cm⁻² is larger than the applied bias (10 mV), in these conditions the output current of the device can be modulated on both side of the zero current level by varying the light intensity. The light modulation of the output current does not affect the spin-polarized charge transport through the device as the photogenerated carriers are non-spin polarized, and the photo-current only shifts the baseline of the magneto current response of the device. Figure 4a shows the output current vs magnetic field (I-B) response of the device under various light irradiation conditions at an applied bias of 10 mV. The overall device current decreases with increasing light intensity. The I-B responses of the device moves from an overall positive current to an overall negative current through a point in which for a particular light intensity the anti-parallel state current can be set to zero. Similarly, for a slightly higher intensity, the parallel state current across the device can be set to zero. This effect could be used to realize a switch having zero power consumption in the standby state. Moreover, at yet another light intensity, the parallel state current can exactly be set equal and opposite of the anti-parallel state current ($I_P = -I_{AP}$) (See Figure S2). This can act as a magnetic current converter. However, in all light intensities the value of $\Delta I = I_P - I_{AP}$ remains the same for the applied bias of 10 mV.

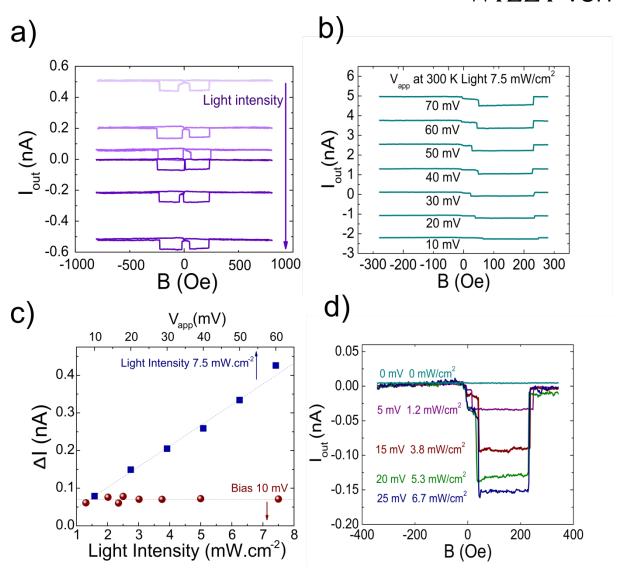


Figure 4. a) Modulation of output current of the device as a function of the magnetic field with varying light intensities at a constant bias of 10 mV. b) I-B responses of the device under constant light irradiation but varying the bias from 10 mV to 70 mV. c) The difference in current in parallel and anti-parallel state (ΔI) as a function of light intensity for a constant bias of 10 mV and as a function of applied bias for a constant light intensity of 7.5 mW.cm⁻². d) Electro-optical modulation of the device under varying applied bias and light intensity to have zero output current in the parallel state, meaning that the generated open-circuit voltage cancels exactly the applied bias.

We now focus on the response for a particular fixed light intensity (in the case shown of 7.5 mW.cm⁻²). Here, the I-B response of the device can be modulated around the zero current level with a modified ΔI under different applied bias (Figure 4b). The modification of ΔI with both light and bias is summarized in Figure 4c. ΔI remains constant with increasing light intensity

for a fixed device bias. This again confirms that the photo-generated carriers do not affect the spin-transport properties through the device. At a fixed bias the amount of spin-polarized charge carriers flowing through the device remains constant whereas for a change in applied bias, this amount changes owing to the change in ΔI . The current level in the parallel state of the device can be tuned both by applied bias and light. The applied bias to the device injects spin-polarized carriers into the molecular layer through molecule/ $Ni_{80}Fe_{20}$ interface whereas the incident light creates non-spin polarized carriers responsible for photocurrent generation. The interplay between light and applied biases can lead to certain states where for an applied bias equal to the open-circuit voltage, the parallel state current can be set to zero. At those particular light intensities and applied biases, spin-polarized current can be generated for the anti-parallel orientation of the electrodes as shown in Figure 4d. This complex electro-optical modulation of the device can lead to the operation of a spin-valve that consumes negligible power in the stand-by state and can act as a spin-photodetector $^{[16,32,33]}$.

In conclusion, we have illustrated the room temperature operation of a spin-photovoltaic device based on a thin film of hole transporting small molecule H₂Pc on a transparent glass substrate. The photovoltaic power conversion efficiency was enhanced as a result of the enhanced photovoltage generation with the application of a small applied magnetic field. The device can act as a perfect magnetic current converter, as a spin-photodetector and can generate purely spin-polarized current under certain applied bias and light irradiation conditions. Our approach opens the door towards the generation of a more complex and highly efficient spin-photovoltaic device, e.g. a p-n heterojunction device, which is a necessary building block for future devices operating at room temperature on transparent substrates.

Experimental Section

Device fabrication:

Vertical spin valves with Co/AlO_x/H₂Pc/Ni₈₀Fe₂₀ have been fabricated *in-situ* in an ultrahigh vacuum (UHV) chamber with base pressure less than 10⁻⁹ mbar. The pyrex glass substrates were cleaned in a ultrasonic bath using acetone and isopropanol subsequently and then dried with a N₂ flow. The substrates were cooled down to liquid N₂ temperature before the deposition of the materials. At first, eight 12 nm thick Co lines were deposited through shadow masking technique to define the bottom electrodes. A 1.5 nm Al layer was then deposited everywhere on the sample and semi-oxidized to form a leaky AlO_x barrier. A 90 nm thick H₂Pc was then deposited again with shadow masking technique to form the molecular layer on Co/AlO_x while leaving one junction without the molecular layer as reference. Finally, the top Ni₈₀Fe₂₀ lines were deposited using another shadow mask to complete the device. The area of the cross-bar geometry devices were ranging from 210×280 μm² to 370×550 μm². Co, Ni₈₀Fe₂₀ and Al were purchased from Lesker (purity: 99.95%) and were used as received. Co, Ni₈₀Fe₂₀ were evaporated from an e-beam evaporator in one of the UHV chamber with a rate of 1 Å/s (for the top Ni₈₀Fe₂₀ deposition, the starting rate was 0.1 Å/s for the first 2 nm to protect the soft organic layer). Al was thermally evaporated from a Knudsen cell at a rate of 1 Å/s. H₂Pc was purchased from Sigma Aldrich (with 99.99% purity) and was used without further purification. The molecules were evaporated at a rate of 0.1 Å/s from another Knudsen cell from a separate chamber.

Thin film characterization and electrical measurements:

Both the metal and the organic materials were calibrated using the quartz crystal monitor during the evaporation. The film thicknesses were measured using X-ray reflectivity (XRR) technique and the morphologies were checked by atomic force microscopy (AFM) technique. The devices were measured in a standard four-probe configuration in a variable temperature Lakeshore probe station (equipped with magnetic field) under high vacuum.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

Acknowledgements

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Supporting Information

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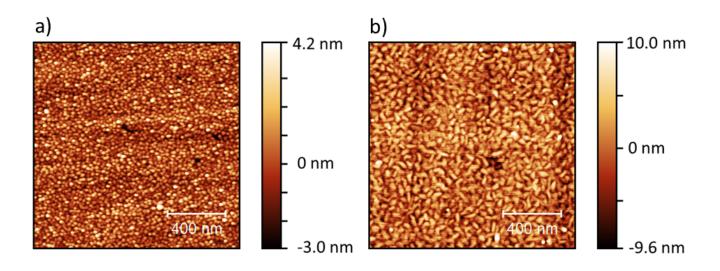


Figure S1. AFM image of 30 nm H₂Pc grown on glass substrate with a) LT condition having rms roughness of 1.9 nm b) RT condition having rms roughness of 2.4 nm

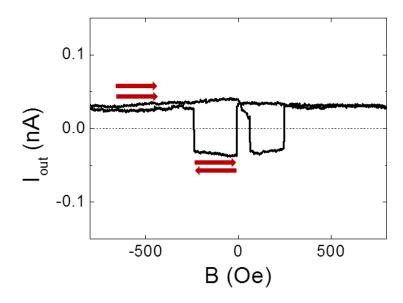


Figure S2. Light modulation of the magnetocurrent leads to a state where the parallel state current is equal and opposite to the anti-parallel state current (device bias: 10 mV; light intensity: 2.5 mW.cm⁻²).

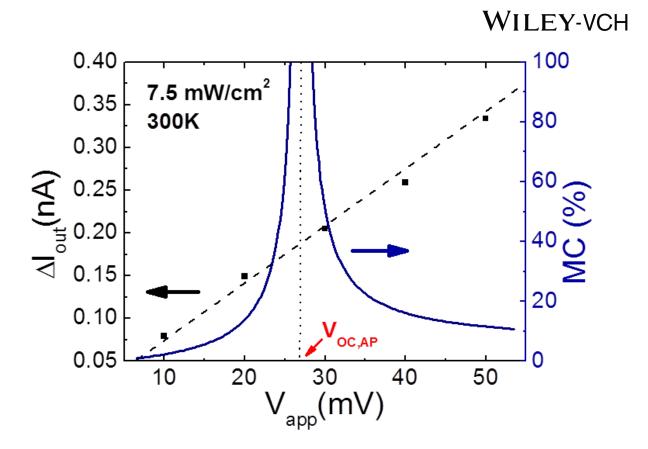


Figure S3. Difference in current in parallel and anti-parallel state (ΔI) and the corresponding MC as a function of the applied bias (V_{app}) (calculated from the linear dependence of ΔI as a function of V_{app}) under light irradiation of 7.5 mW.cm⁻². At an applied bias equal to the open circuit voltage, the MC value tends to infinite making it favorable for a light sensor.